# Non-Standard Interactions: Current Status & Future Prospects



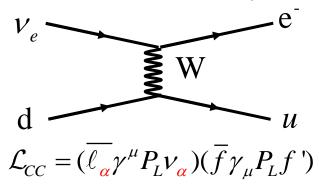


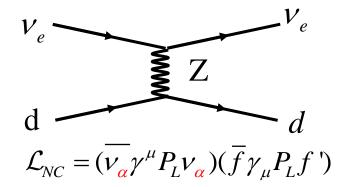
### Outline

- Non-standard neutrino interactions (NSIs)
- Neutrino oscillations with NSIs
- Current status
- Future prospects
- Probing NSIs at the ESSvSB experiment
- Conclusions (and comments)

### Non-standard neutrino interactions: NSIs

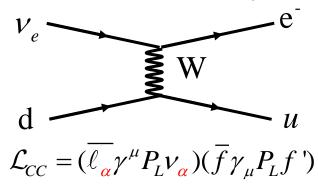
In the Standard Model,

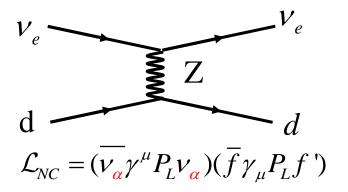




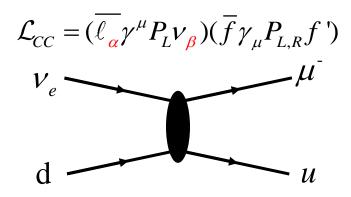
### Non-standard neutrino interactions: NSIs

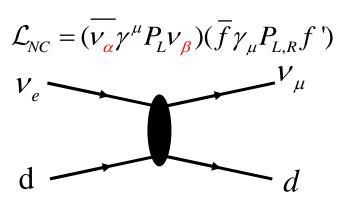
In the Standard Model,





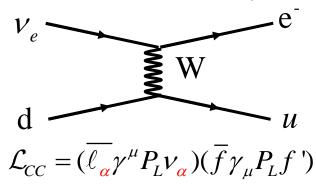
With new physics, we could have

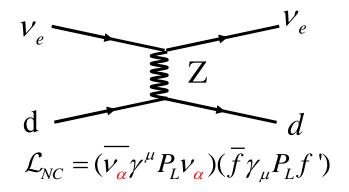




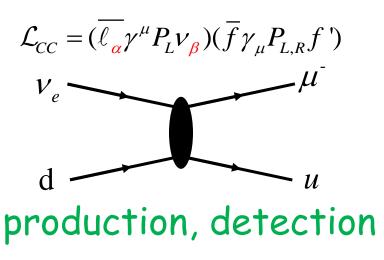
### Non-standard neutrino interactions: NSIs

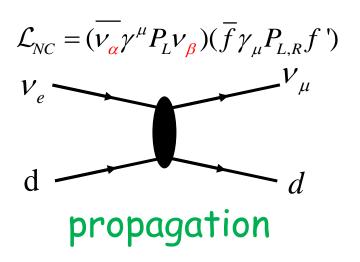
In the Standard Model,





With new physics, we could have





### Neutrino oscillations with NSIs

Neutrino states at the source and detector:

$$|v_{\alpha}^{s}\rangle = |v_{\alpha}\rangle + \sum_{\gamma=e,\mu,\tau} \varepsilon_{\alpha\gamma}^{s} |v_{\gamma}\rangle \qquad ; \qquad \langle v_{\beta}^{d}| = \langle v_{\beta}| + \sum_{\gamma=e,\mu,\tau} \varepsilon_{\beta\gamma}^{d} \langle v_{\gamma}|$$

Matter effects in the presence of NSIs:

$$A^{NSI} = egin{pmatrix} A & 0 & 0 \ 0 & 0 & 0 \ 0 & 0 & 0 \end{pmatrix} + A egin{pmatrix} arepsilon_{ee}^m & arepsilon_{e\mu}^m & arepsilon_{e\mu}^m & arepsilon_{\mu\mu}^m & arepsilon_{\mu\tau}^m \ arepsilon_{ au e}^m & arepsilon_{ au\mu}^m & arepsilon_{ au au}^m \end{pmatrix}$$

 9x2 (source) + 9x2 (detector) + 9 (propagation) new NSI parameters

### Neutrino oscillations with NSIs

Neutrino states at the source and detector:

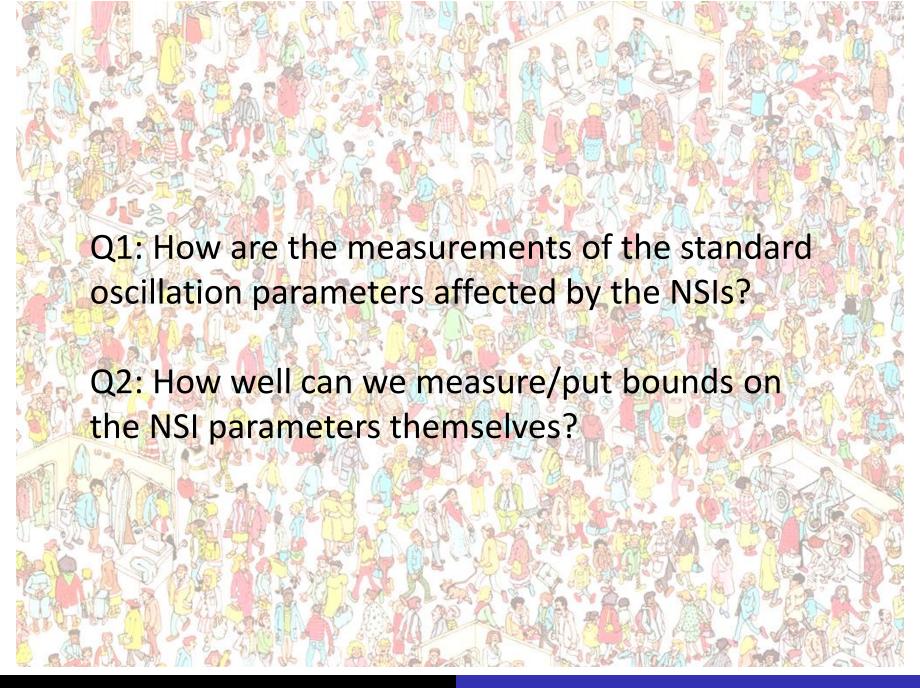
$$|v_{\alpha}^{s}\rangle = |v_{\alpha}\rangle + \sum_{\gamma=e,\mu,\tau} \varepsilon_{\alpha\gamma}^{s} |v_{\gamma}\rangle \qquad ; \qquad \langle v_{\beta}^{d}| = \langle v_{\beta}| + \sum_{\gamma=e,\mu,\tau} \varepsilon_{\beta\gamma}^{d} \langle v_{\gamma}|$$

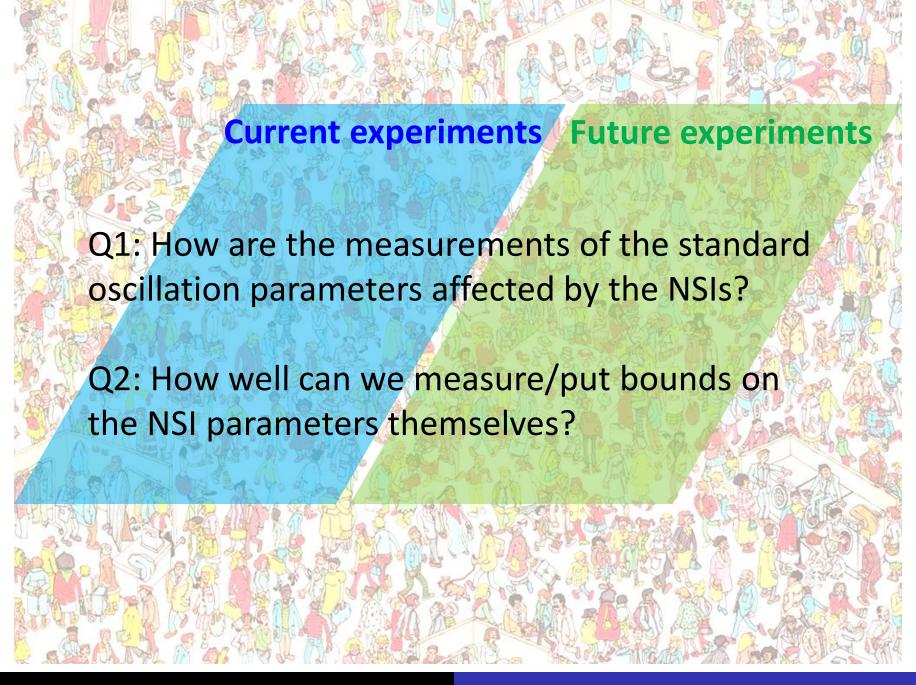
Matter effects in the presence of NSIs:

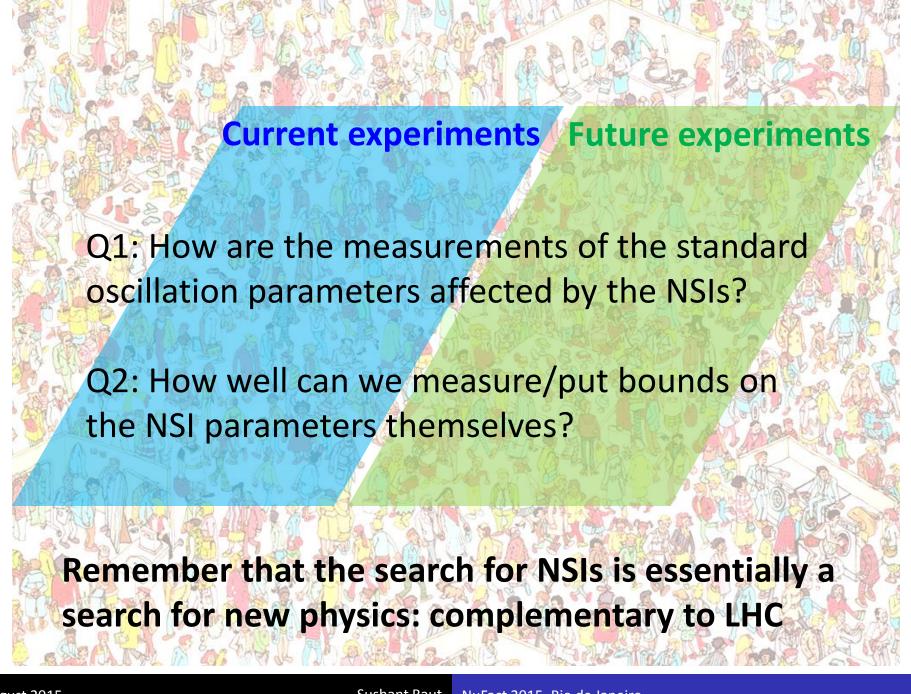
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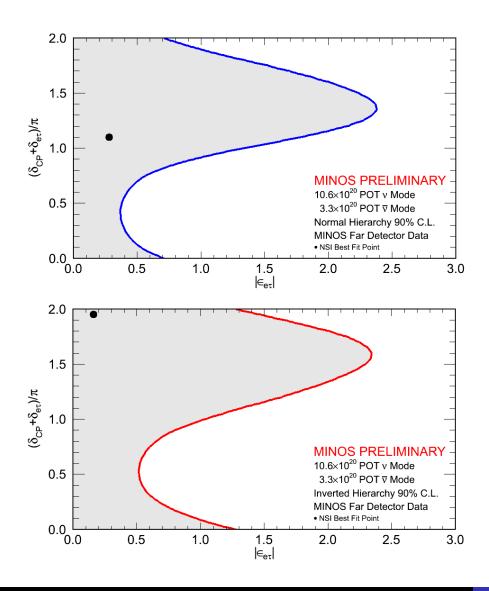
9x2 (source) + 9x2 (detector) + 9 (propagation) new NSI parameters ... a slight expansion of the parameter space!





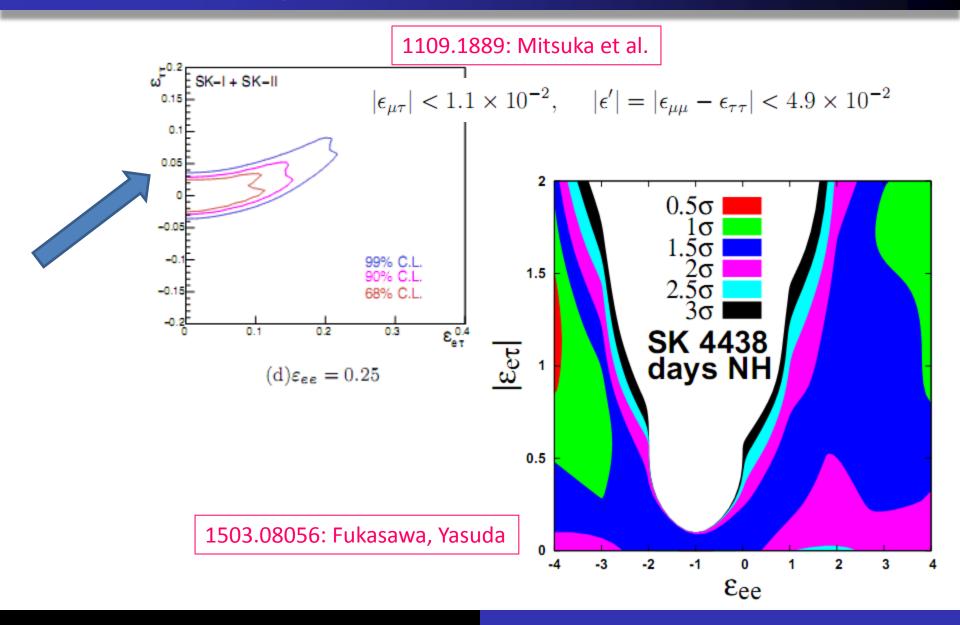


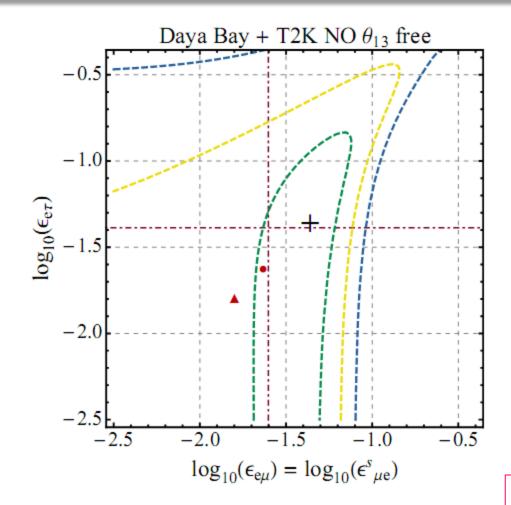




#### MINOS beam data

Remember: Stringent constraints from the charged lepton sector!





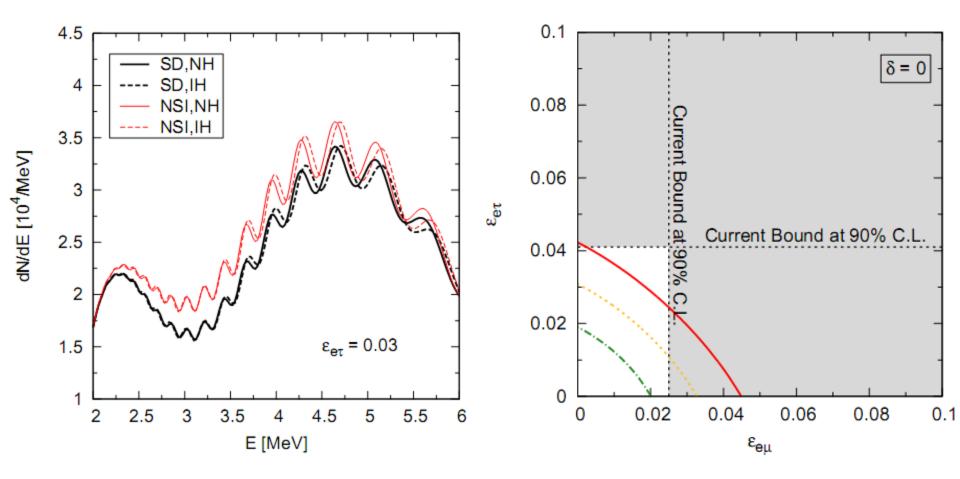
Discrepancy between  $\theta_{13}$  value measured at T2K and at the reactor experiments: role of NSIs

1405.0416: Girardi, Meloni, Petcov

- CKM unitarity
- Pion decays
- Zero distance oscillations at KARMEN, NOMAD
- Loop bounds

$$|\varepsilon_{\alpha\beta}^{ud}| < \begin{pmatrix} 0.041 & 0.025 & 0.041 \\ 1.8 \cdot 10^{-6} & 0.078 & 0.013 \\ 0.026 & 0.013 & 0.013 \\ 0.12 & 0.018 & 0.13 \end{pmatrix}$$

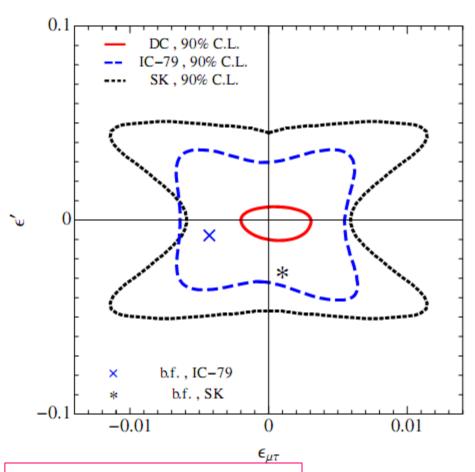
0907.0097: Biggio, Blennow, Fernandez-Martinez



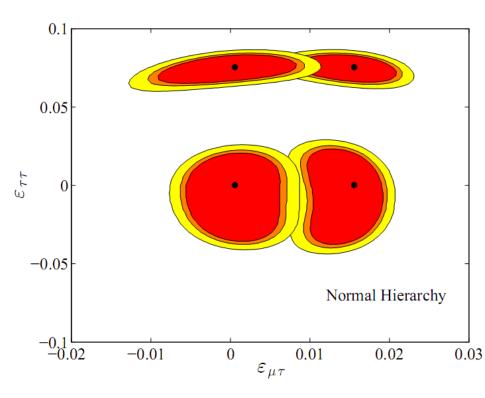
Medium baseline reactors: JUNO

1310.5917: Ohlsson, Zhang, Zhou

### IceCube & Deep Core

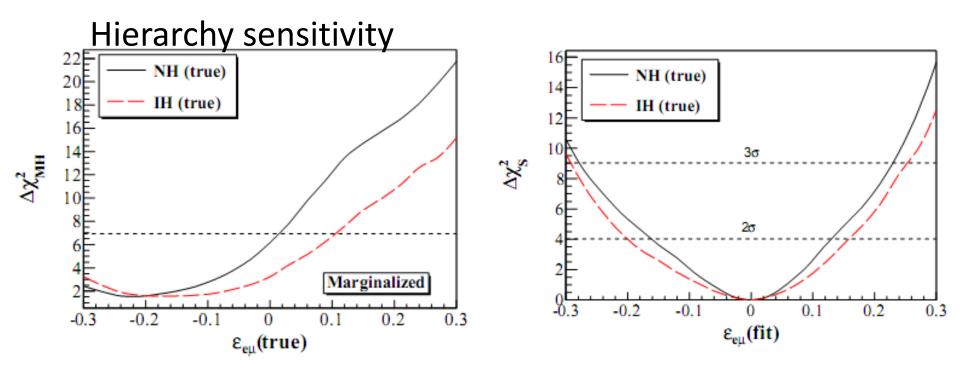


#### **PINGU**



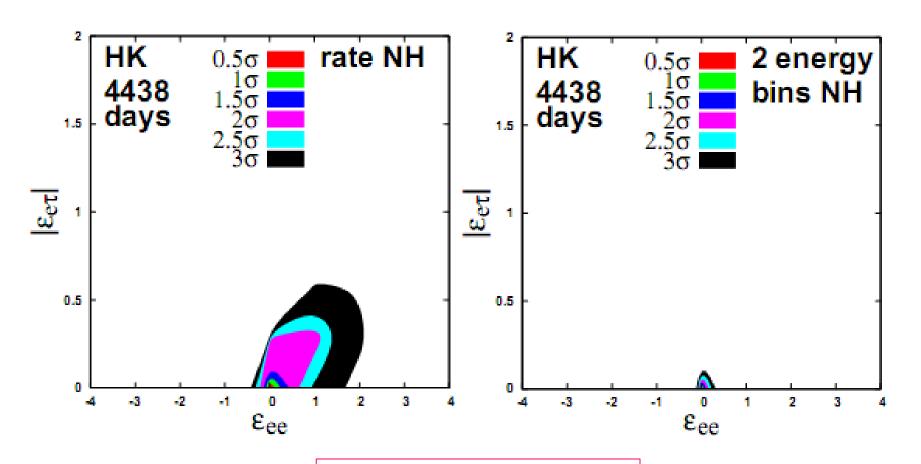
1410.0410: Choubey, Ohlsson

1304.1042: Esmaili, Smirnov



ICAL @ INO

1507.02211: Choubey, Ghosh, Ohlsson, Tiwari

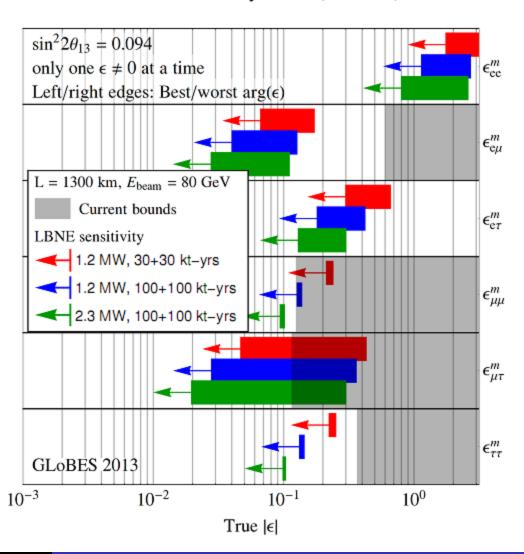


1503.08056: Fukasawa, Yasuda

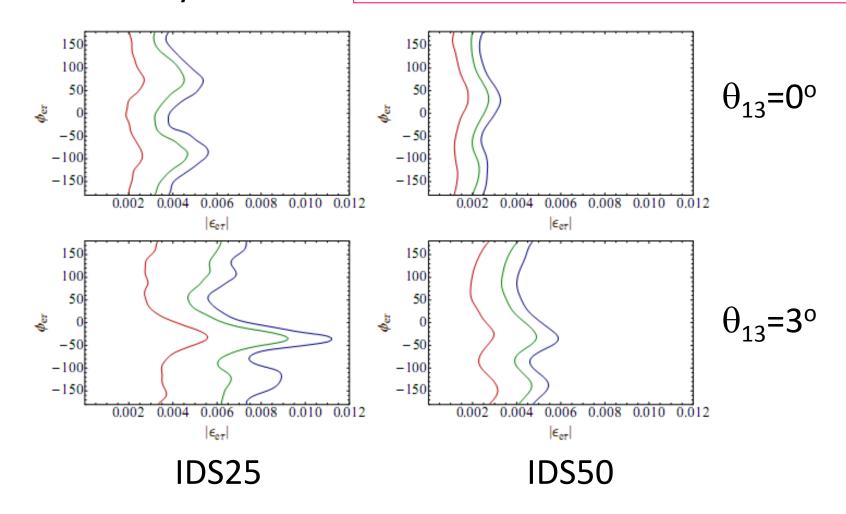
### LBNE/LBNF/DUNE

1307.7335: Adams et al.

#### NC NSI discovery reach ( $3\sigma$ C.L.)



Neutrino Factory: 4000 km | 1105.5936: Coloma, Donini, Lopez-Pavon, Minakata



## Probing NSIs at ESSvSB

1507.02868: Blennow, Choubey, Ohlsson, SR



## The proposed ESSvSB experiment



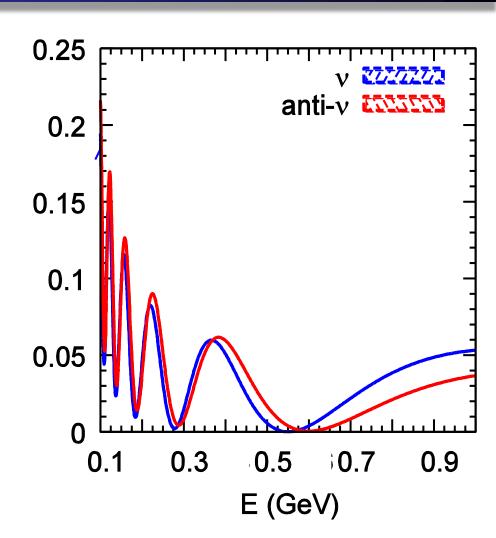
- European Spallation Source (ESS): under construction in Lund, Sweden
- Proposal to use the proton beam to produce a beam of neutrinos – peak energy 250 MeV

- Possible site for detector: mine in Garpenberg, Sweden – 540 km
- The mine can host a MEMPHYS-like Water Cerenkov detector

1309.7022: Baussan et al.

### **ESSVSB**

- For a 540 km baseline, the second oscillation maximum (which is sensitive to  $\delta$ ) is at 400 MeV
- The peak energy of the ESSvSB unoscillated spectrum lies at this energy, giving this experiment good sensitivity to  $\delta$
- How will this sensitivity be affected by NSIs? Can we use ESSvSB to measure NSIs?



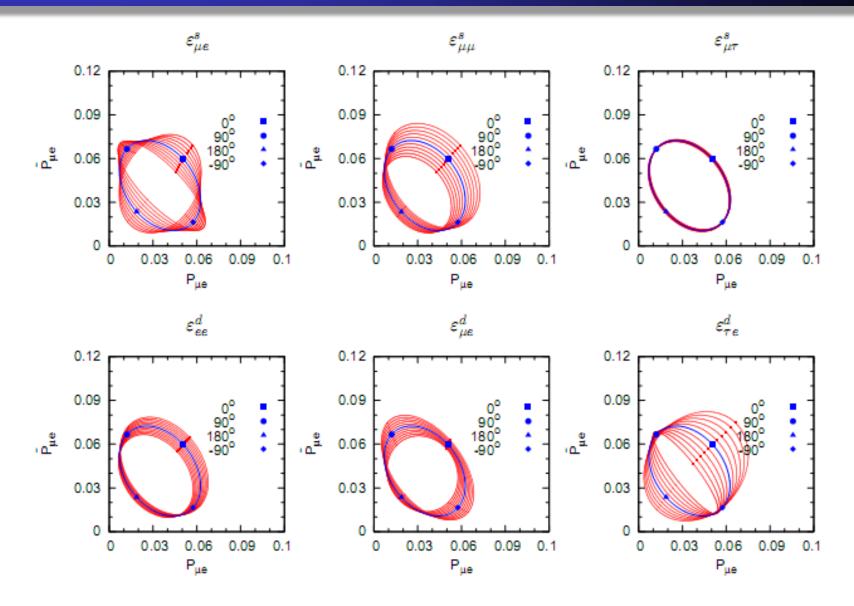
### Neutrino oscillations with NSIs

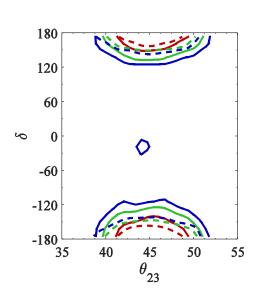
 Propagation NSIs are only relevant for large matter effects and high energy. Therefore we only consider source and detector NSIs here.

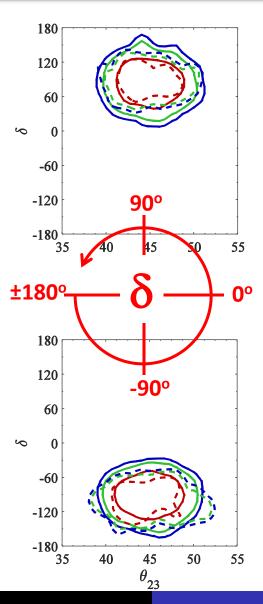
$$|v_{\alpha}^{s}\rangle = |v_{\alpha}\rangle + \sum_{\gamma=e,\mu,\tau} \varepsilon_{\alpha\gamma}^{s} |v_{\gamma}\rangle \qquad ; \qquad \langle v_{\beta}^{d}| = \langle v_{\beta}| + \sum_{\gamma=e,\mu,\tau} \varepsilon_{\beta\gamma}^{d} \langle v_{\gamma}|$$

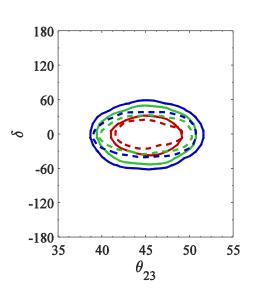
• Since ESSvSB will only observe  $v_{\mu} \rightarrow v_{\mu}$  and  $v_{\mu} \rightarrow v_{e}$  channels, the relevant NSI parameters are

## Interplay of NSI parameters with $\delta$





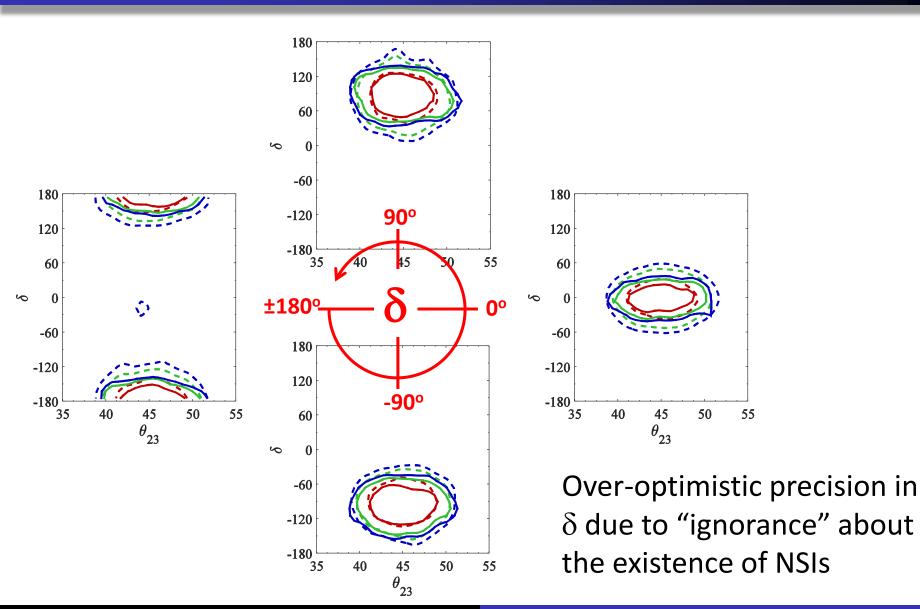


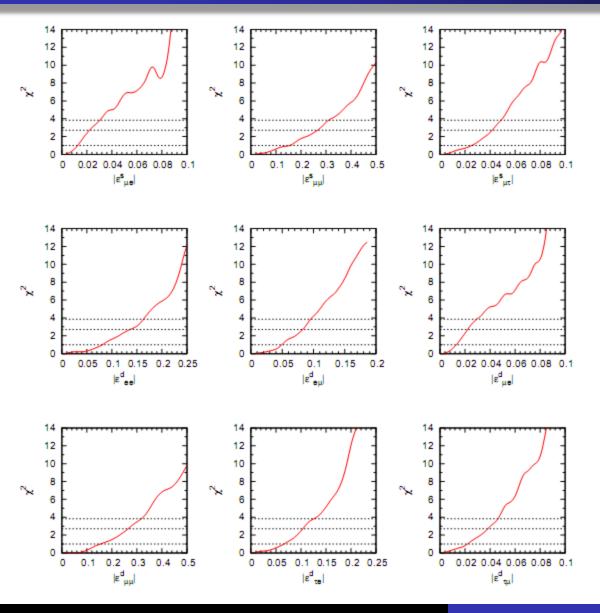


Worsening of  $\delta$ -precision because of NSIs

### Salient features:

- Same features seen across all  $\theta_{23}$
- $\theta_{23}$ -precision is largely unaffected by NSIs
- Worsening of  $\delta$  -precision is most for  $\delta$  =180
- Worsening of precision is at most twice as bad: The measurement of  $\delta$  at ESSvSB is quite robust against NSIs





Limits on NSI parameters from ESSvSB

Parameter	Realistic limits	Optimistic limits	Existing limits
$arepsilon_{ m \mu e}^{ m s}$	0.025	0.014	0.026
$\epsilon_{~\mu\mu}^{s}$	0.27	0.27	0.078
${f \epsilon}^{\sf s}_{\ \mu au}$	0.040	0.040	0.013
$\epsilon_{ee}^{d}$	0.15	0.15	0.041
$\epsilon_{\rm e\mu}^{\rm d}$	0.087	0.082	0.026
$arepsilon_{ m \mu e}^{ m d}$	0.025	0.014	0.025
٤ <sup>d</sup>	0.28	0.27	0.078
$\epsilon_{ au_{ au e}}^{ au}$	0.11	0.12	0.041
$\epsilon_{ au\mu}^{d}$	0.040	0.033	0.013

### Conclusions & comments

- Non-standard effects arise from various possible BSM scenarios, and can affect neutrino oscillations.
- Current bounds on NSI parameters are at the 10<sup>-2</sup> level. Future experiments like PINGU and HK will impose more stringent bounds. The neutrino factory can push them even further.
- We need more studies on source/detector NSIs.
- We need studies in the full parameter space, unless it is constrained by physical reasons. Yes, that includes the phases!
- The measurement of  $\delta$  at ESSvSB is robust to the presence of NSIs
- ESSvSB can impose strong bounds on  $\epsilon^{s}_{\mu e}$  and  $\epsilon^{d}_{\mu e}$ , which are more restrictive than the current bounds



# Analytical formula

$$\begin{split} P_{\nu_{\mu}^{s} \to \nu_{e}^{d}}^{\text{vac}} &= 4s_{13}^{2}s_{23}^{2} \sin^{2} \frac{\Delta m_{31}^{2}L}{4E} \\ &+ \left(\frac{\Delta m_{21}^{2}}{\Delta m_{31}^{2}}\right)^{2} c_{23}^{2}s_{2\times 12}^{2} \left(\frac{\Delta m_{31}^{2}L}{4E}\right)^{2} \\ &+ \frac{\Delta m_{21}^{2}}{\Delta m_{31}^{2}} s_{13} s_{2\times 12} s_{2\times 23} \cos \delta_{\text{CP}} \frac{\Delta m_{31}^{2}L}{4E} \sin \frac{\Delta m_{31}^{2}L}{2E} \\ &- 2\frac{\Delta m_{21}^{2}}{\Delta m_{31}^{2}} s_{13} s_{2\times 12} s_{2\times 23} \sin \delta_{\text{CP}} \frac{\Delta m_{31}^{2}L}{4E} \sin^{2} \frac{\Delta m_{31}^{2}L}{4E} \\ &- 4|\epsilon_{\mu e}^{s}| s_{13} s_{23} \cos(\phi_{\mu e}^{s} + \delta_{\text{CP}}) \sin^{2} \frac{\Delta m_{31}^{2}L}{4E} \\ &- 2|\epsilon_{\mu e}^{s}| s_{13} s_{23} \sin(\phi_{\mu e}^{s} + \delta_{\text{CP}}) \sin \frac{\Delta m_{31}^{2}L}{2E} \\ &- 4|\epsilon_{\mu e}^{d}| s_{13} c_{2\times 23} s_{23} \cos(\phi_{\mu e}^{d} + \delta_{\text{CP}}) \sin^{2} \frac{\Delta m_{31}^{2}L}{4E} \\ &- 2|\epsilon_{\mu e}^{d}| s_{13} s_{23} \sin(\phi_{\mu e}^{d} + \delta_{\text{CP}}) \sin \frac{\Delta m_{31}^{2}L}{4E} \\ &- 2|\epsilon_{\mu e}^{d}| s_{13} s_{23} \sin(\phi_{\mu e}^{d} + \delta_{\text{CP}}) \sin \frac{\Delta m_{31}^{2}L}{4E} \end{split}$$

0708.0152: Kopp, Lindner, Ota, Sato

$$+ 4|\epsilon_{\tau e}^{d}|s_{13}s_{2\times23}s_{23}\cos(\phi_{\tau e}^{d} + \delta_{\text{CP}})\sin^{2}\frac{\Delta m_{31}^{2}L}{4E}$$

$$-|\epsilon_{\mu e}^{s}|\frac{\Delta m_{21}^{2}}{\Delta m_{31}^{2}}s_{2\times12}c_{23}\sin\phi_{\mu e}^{s}\frac{\Delta m_{31}^{2}L}{2E}$$

$$+2|\epsilon_{\mu e}^{d}|\frac{\Delta m_{21}^{2}}{\Delta m_{31}^{2}}s_{2\times12}s_{23}^{2}c_{23}\cos\phi_{\mu e}^{d}\frac{\Delta m_{31}^{2}L}{4E}\sin\frac{\Delta m_{31}^{2}L}{2E}$$

$$-|\epsilon_{\mu e}^{d}|\frac{\Delta m_{21}^{2}}{\Delta m_{31}^{2}}s_{2\times12}c_{23}\sin\phi_{\mu e}^{d}\frac{\Delta m_{31}^{2}L}{2E}$$

$$\cdot \left[1-2s_{23}^{2}\sin^{2}\frac{\Delta m_{31}^{2}L}{2E}\right]$$

$$+2|\epsilon_{\tau e}^{d}|\frac{\Delta m_{21}^{2}}{\Delta m_{31}^{2}}s_{2\times12}s_{23}c_{23}^{2}\cos\phi_{\mu e}^{d}\frac{\Delta m_{31}^{2}L}{4E}\sin\frac{\Delta m_{31}^{2}L}{2E}$$

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$$+2|\epsilon_{\tau e}^{d}|\frac{\Delta m_{21}^{2}}{\Delta m_{31}^{2}}s_{2\times12}s_{23}c_{23}^{2}\sin\phi_{\mu e}^{d}\frac{\Delta m_{31}^{2}L}{2E}\sin^{2}\frac{\Delta m_{31}^{2}L}{4E}$$

$$+\mathcal{O}\left(\left[\frac{\Delta m_{21}^{2}}{\Delta m_{21}^{2}}\right]^{3}\right)+\mathcal{O}\left(\left[\frac{\Delta m_{21}^{2}}{\Delta m_{21}^{2}}\right]^{2}s_{13}\right)+\mathcal{O}\left(\frac{\Delta m_{21}^{2}}{\Delta m_{21}^{2}}s_{13}^{2}\right)$$

### **Current status**

$$\begin{split} |\varepsilon_{\mu e}^{s}| &< 0.026 \qquad , \qquad |\varepsilon_{\mu \mu}^{s}| &< 0.078 \qquad , \qquad |\varepsilon_{\mu \tau}^{s}| &< 0.013 \\ |\varepsilon_{e e}^{d}| &< 0.041 \qquad , \qquad |\varepsilon_{\mu e}^{d}| &< 0.025 \qquad , \qquad |\varepsilon_{\tau e}^{d}| &< 0.041 \\ |\varepsilon_{e \mu}^{d}| &< 0.026 \qquad , \qquad |\varepsilon_{\mu \mu}^{d}| &< 0.078 \qquad , \qquad |\varepsilon_{\tau \mu}^{d}| &< 0.013 \end{split}$$

0907.0097: Biggio, Blennow, Fernandez-Martinez